

Nonadiabatic quantum state manipulation of superconducting structures

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FIAN / NEC / RIKEN

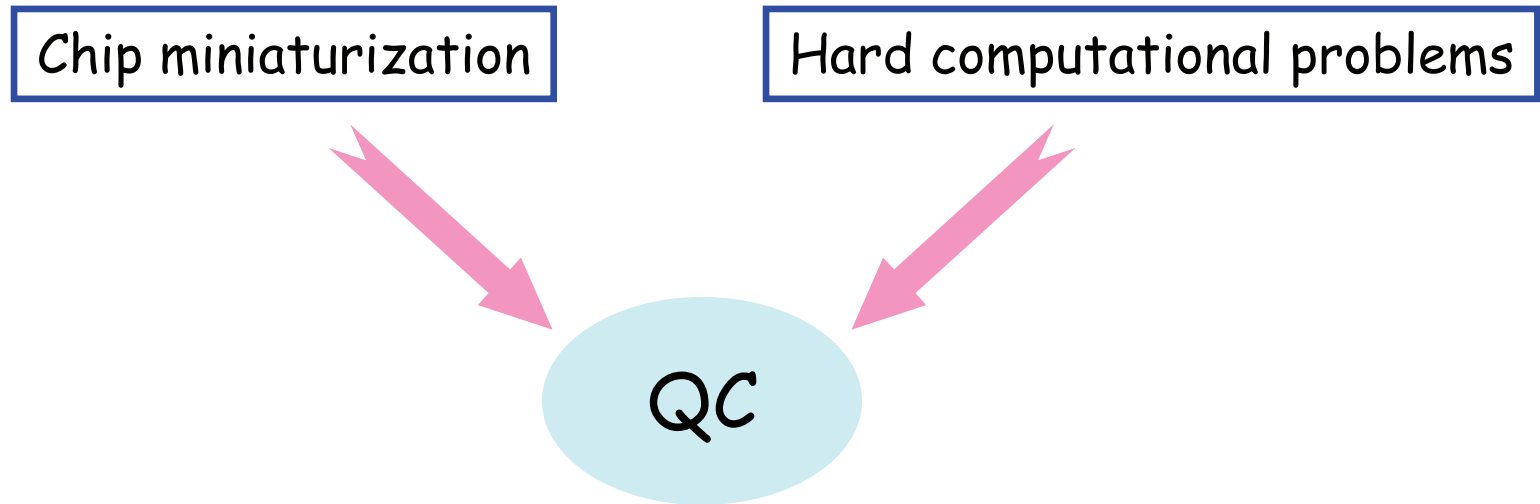
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Jaw-Shen Tsai		

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Why quantum computers?



Shrinking electronics

Electronics have been getting smaller and faster

- Moore's law: computer performance doubles every 18 months
- atomic levels will be reached in ?? Years

R. Feynmann (1985):

"it seems that the laws of physics present no barrier to reducing the size of computers until bits are the size of atoms, and quantum behavior holds sway"

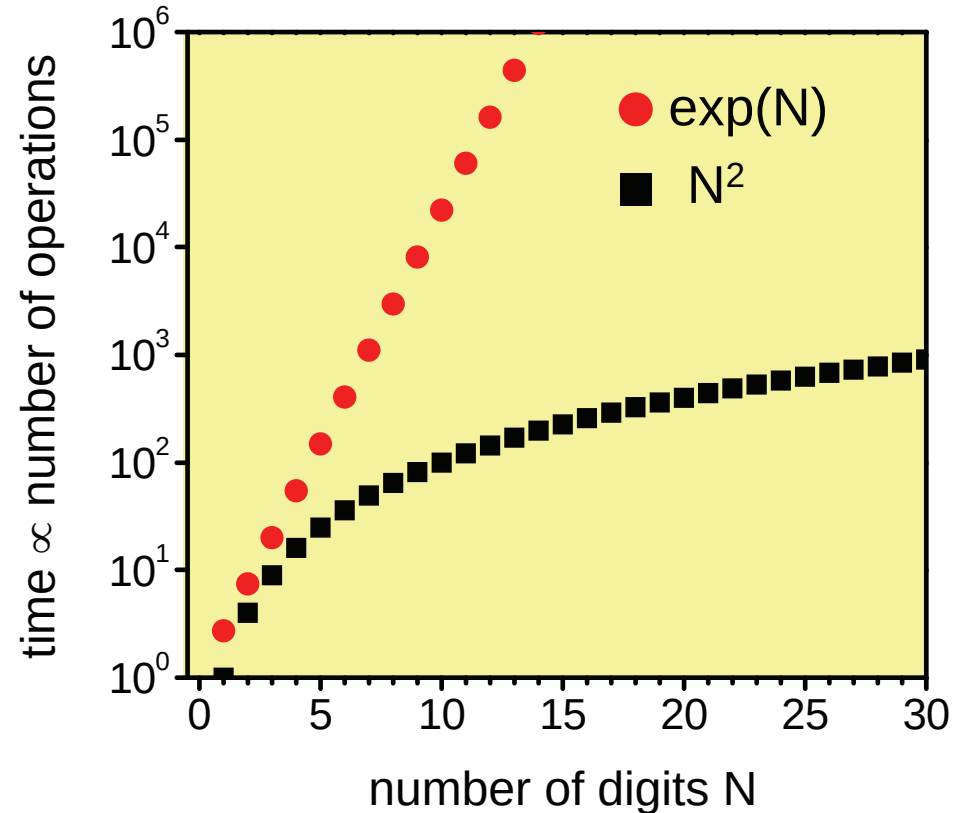
Complexity of computational problems

Complexity of a problem can be:
(N = number of digits in input)

polynomial $N_{op} \propto N^a$

non-polynomial $N_{op} \propto \exp(N)$

Computational time is proportional
to the number of operations



Ex: Factorization of an integer in prime factors is NP problem
We know how to solve the problem but we do not have time!

Power of QC

Factorization of large integers (N digits)

Classical algorithms:

N = 129	8 months (1994)
N = 250	10^6 years (estimated)
N = 1000	10^{25} years (estimated)

Quantum algorithm (P. Shor, 1994):

N = 1000	a few seconds
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Hard computational problems

- factorization of large numbers --> *Shor's algorithm*
- database search --> *Grover's algorithm*
- graph isomorphism
- travelling salesman
- tailor's problem

Main features of QC

qu-bits: $\begin{array}{|c|} \hline 1 \text{ ---} \\ \hline 0 \text{ ---} \\ \hline \end{array} \quad \left| \begin{array}{l} |\psi\rangle = \alpha|0\rangle + \beta|1\rangle \\ |\alpha|^2 + |\beta|^2 = 1 \end{array} \right| \begin{array}{l} |0\rangle \text{ with probability } |\alpha|^2 \\ |1\rangle \text{ with probability } |\beta|^2 \end{array}$

qu-registers: a set of qubits, N qubits $\rightarrow 2^N$ states

qu-gates: input - superposition of states

N - bit register represents 2^N states $\rightarrow$ quantum parallelism

Any operation is performed on all the states at the same time $\rightarrow$ exponential amount of computational space with linear amount of physical space

To build a QC we need: (DiVincenzo criteria)

scalable qubits = two-level quantum systems

initialization = prepare qubit in a given state

state manipulation

read-out

quantum logic gates

"long" coherence time ($\tau \gg 1/f_{\text{clock}}$)

Two-level systems

non solid-state

- atoms
- ions
- nuclear spins
- photons
- ...

solid-state

- CP box ← this work
- RF SQUID
- 3J SQUID
- single J junction
- quantum dots
- ...

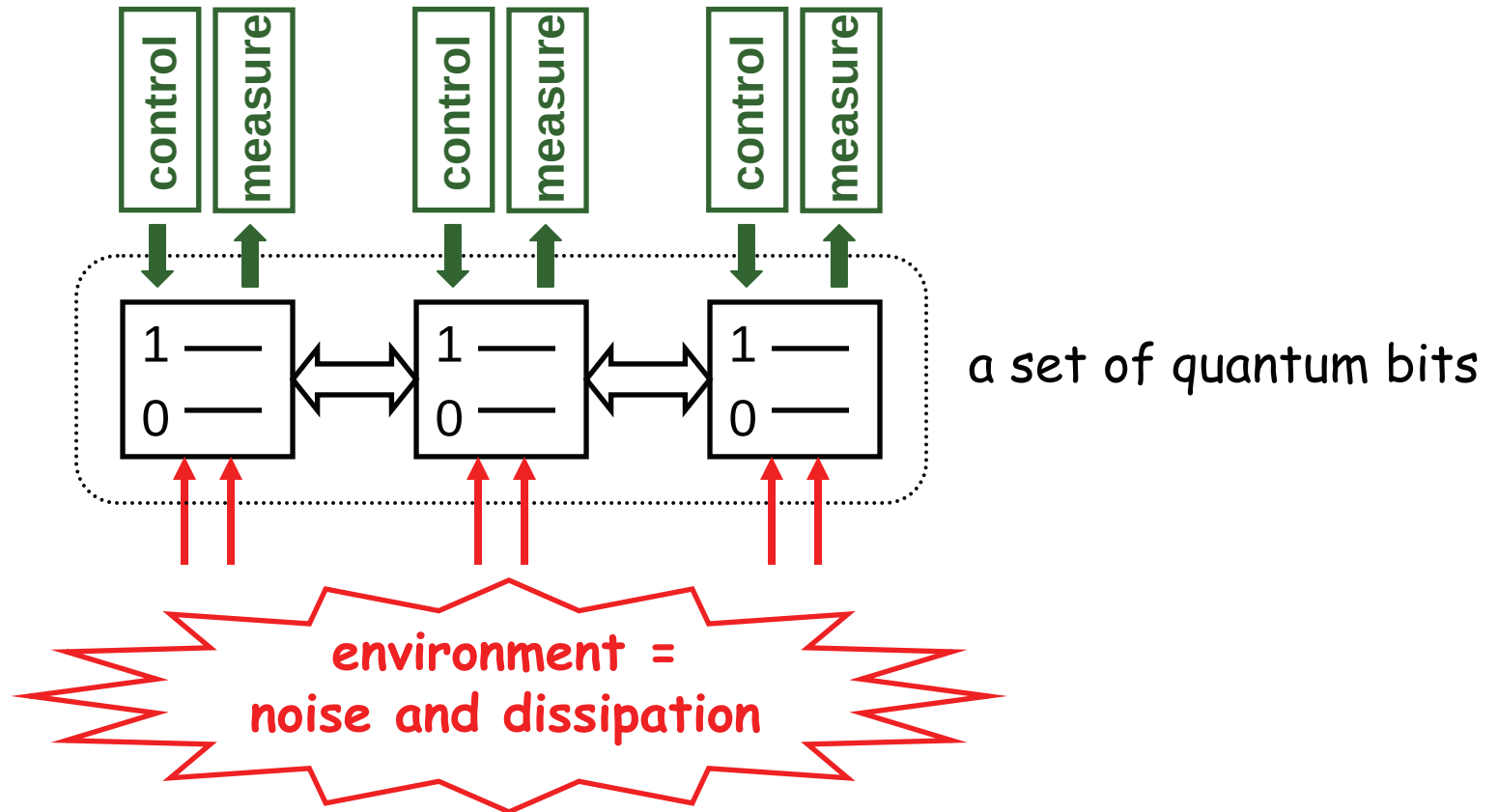
coherence easier



coupling easier



Challenge: coupling vs. decoherence



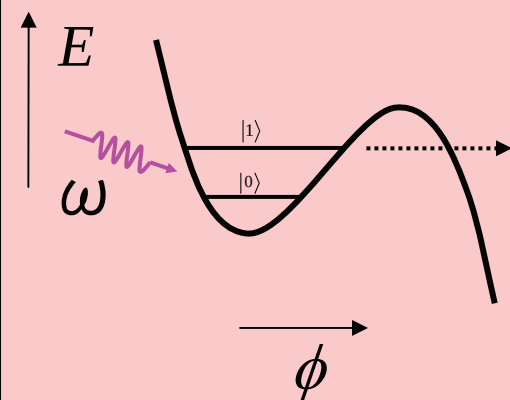
Experimental challenge:
couple qubits to each other, control, & measure
not noise and dissipation

Josephson-junction-based qubits

Single junction

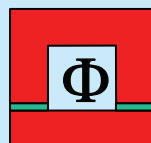


$$U_J = E_J (1 - \cos \phi)$$

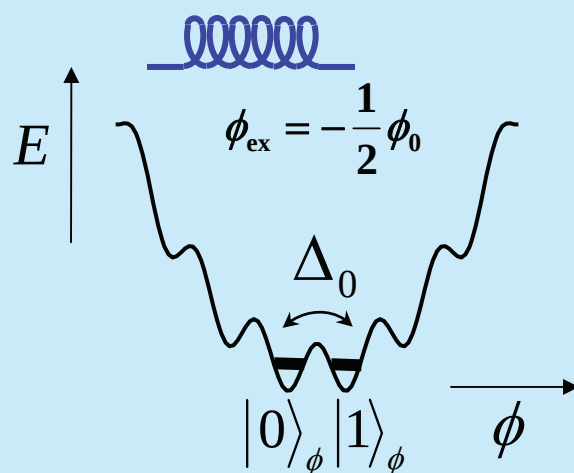


NIST
Kansas
Maryland

SQUID

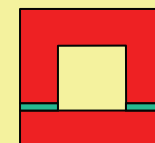


$$E_c < E_J$$

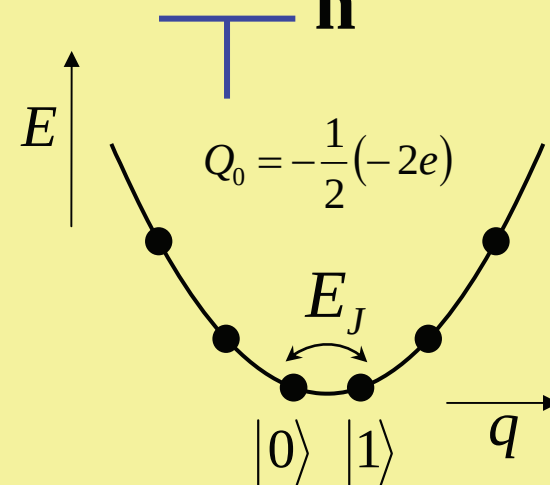


Delft
NTT
Jena

Cooper-pair box



$$E_c > E_J$$



Saclay

NEC
Chalmers
Yale
JPL

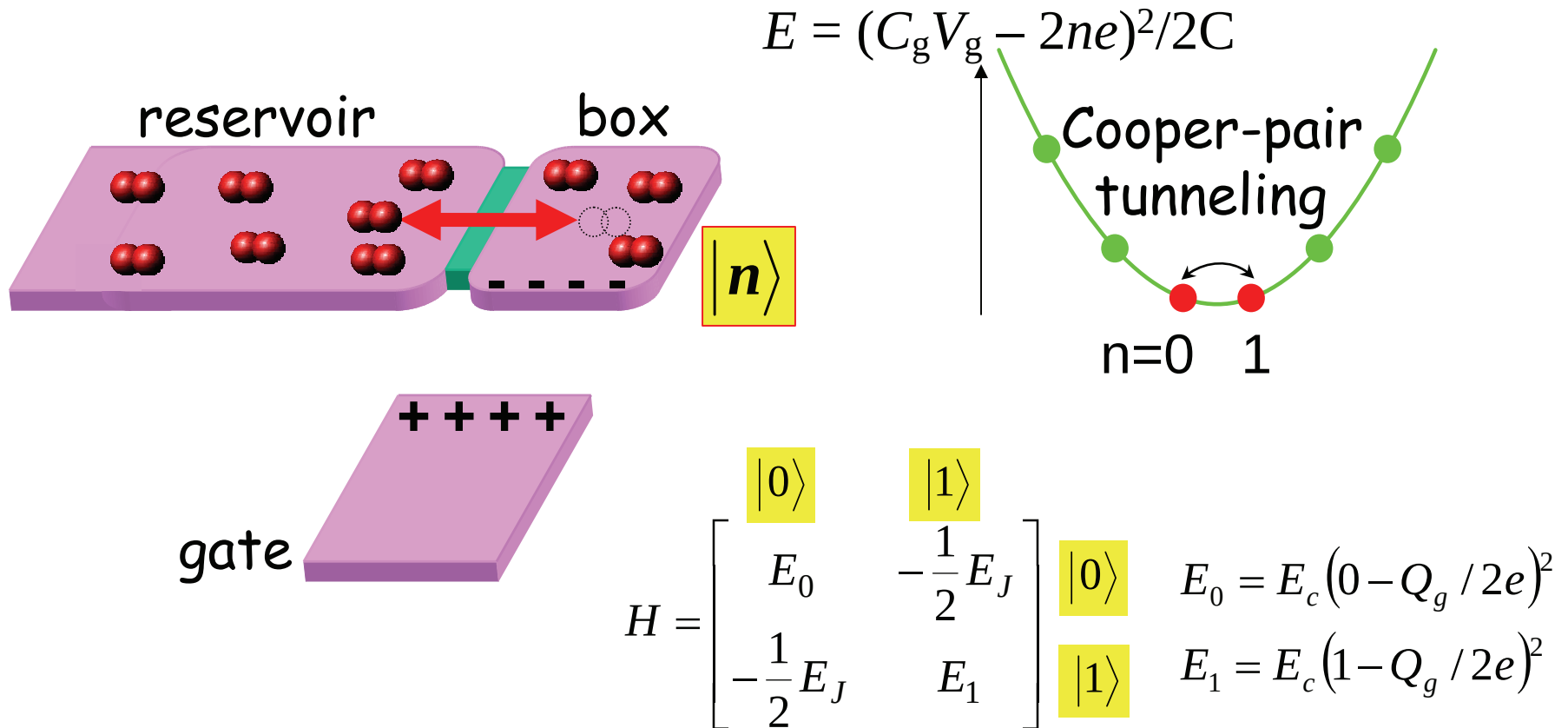
Cooper-pair box

M. Büttiker, 1987
V. Bouchiat et al, 1995

- a **single** artificial two-level system
- $\sim 10^8$ conduction electrons in the box

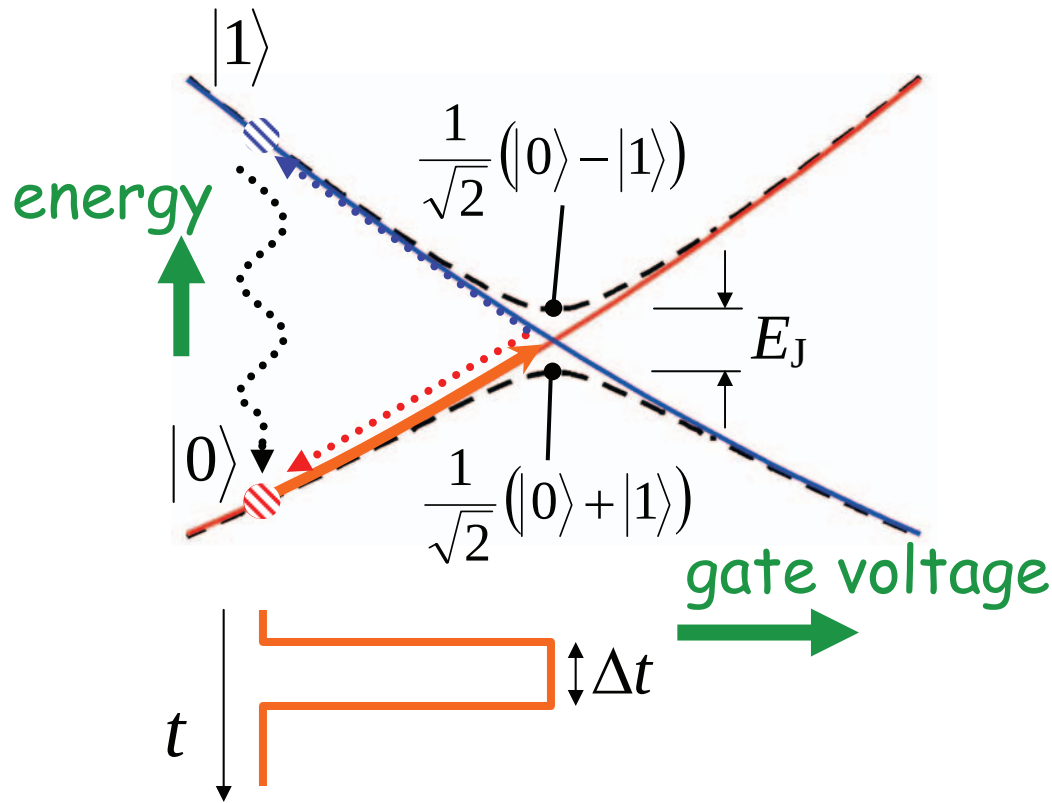
$$\Delta > E_C$$

$$4E_C > E_J \gg kT$$



Charge qubit based on Cooper-pair box

Y. Nakamura et al, 1999



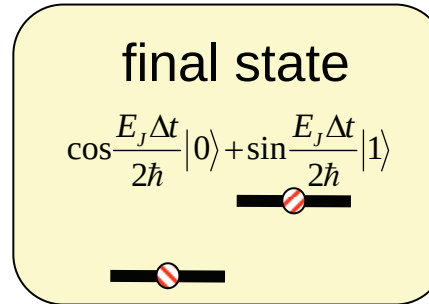
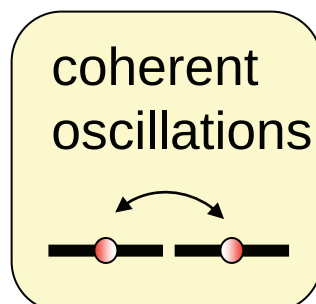
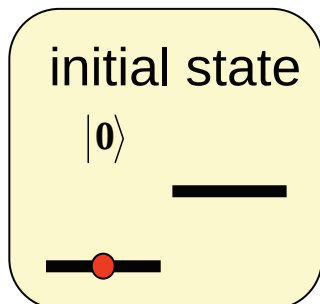
eigenstates: $|\psi_G\rangle, |\psi_E\rangle$

$$\Delta E(Q_0) = \sqrt{\delta E(Q_0)^2 + E_J^2}$$

charge states: $|0\rangle, |1\rangle$

$$\delta E(Q_0) \equiv E_c(Q_0 / e - 1)$$

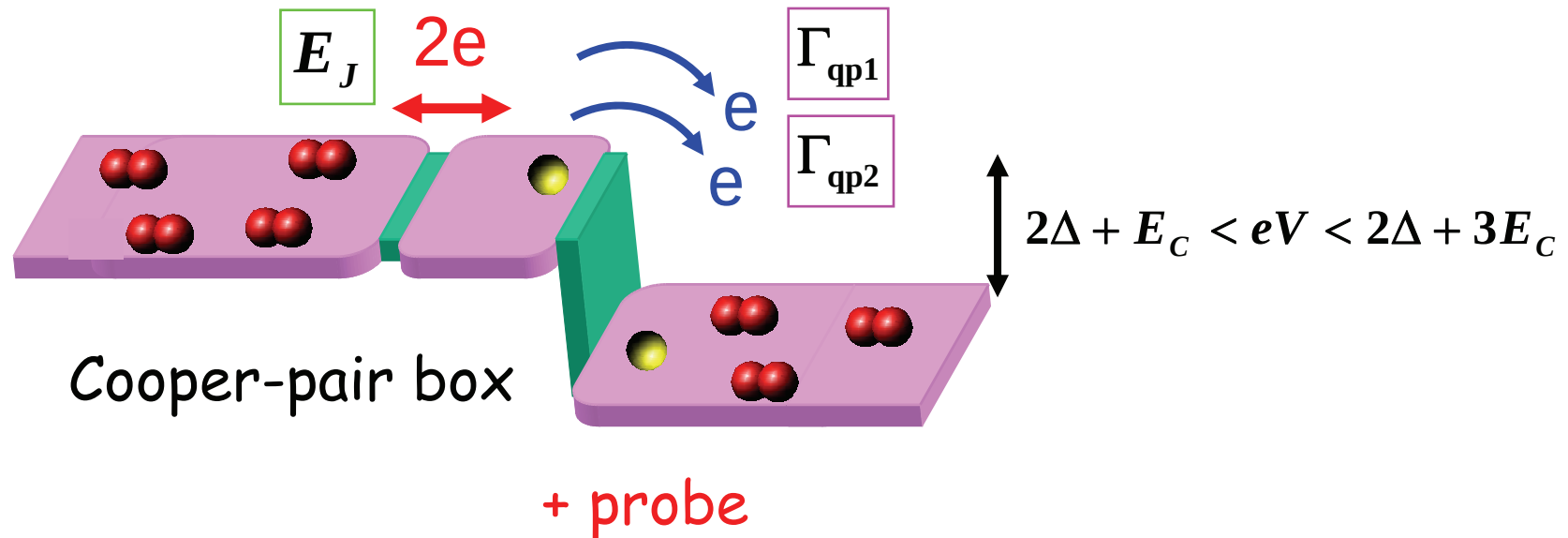
initialization
coherent superposition
read-out



$$p(1) = \frac{1}{2} \left[1 - \cos \left(\frac{E_J \Delta t}{\hbar} \right) \right]$$

Final state read-out

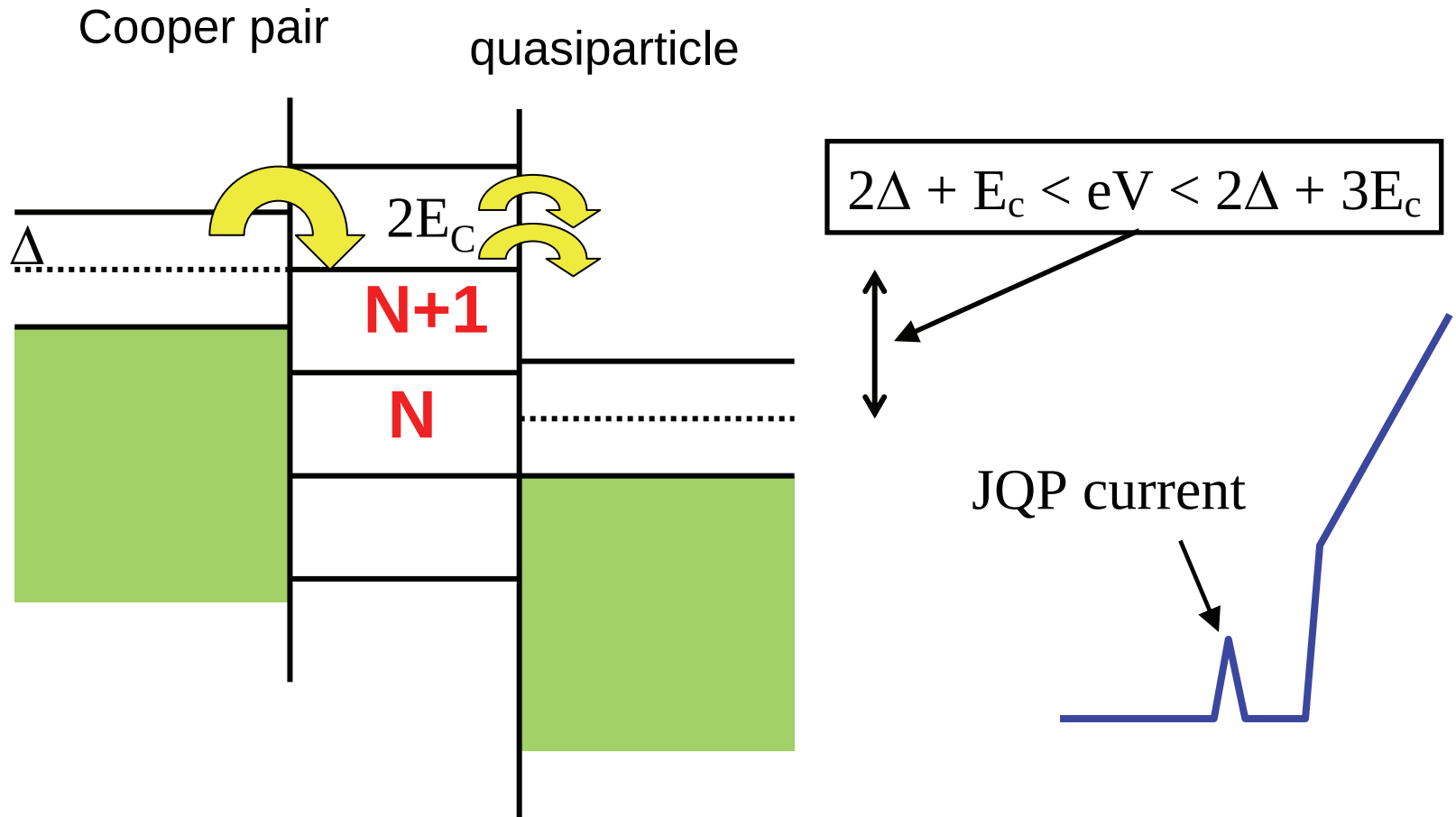
Josephson-quasiparticle cycle (Fulton et al., 1989)



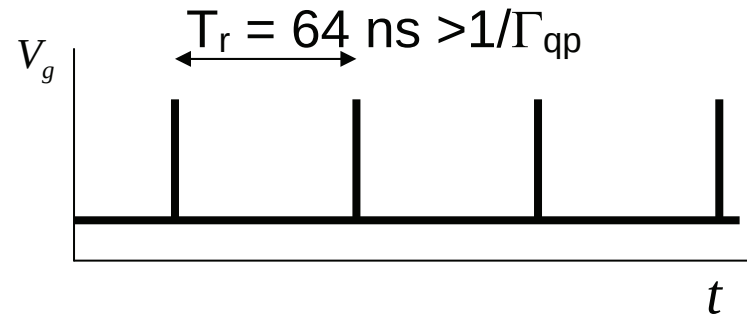
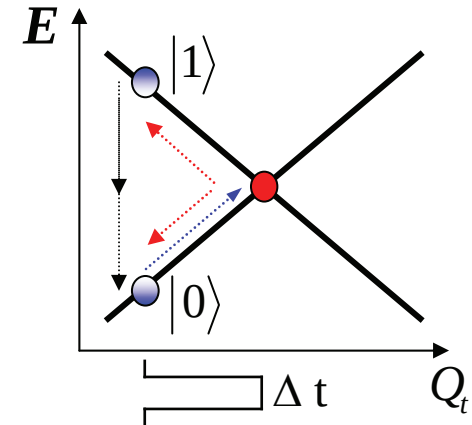
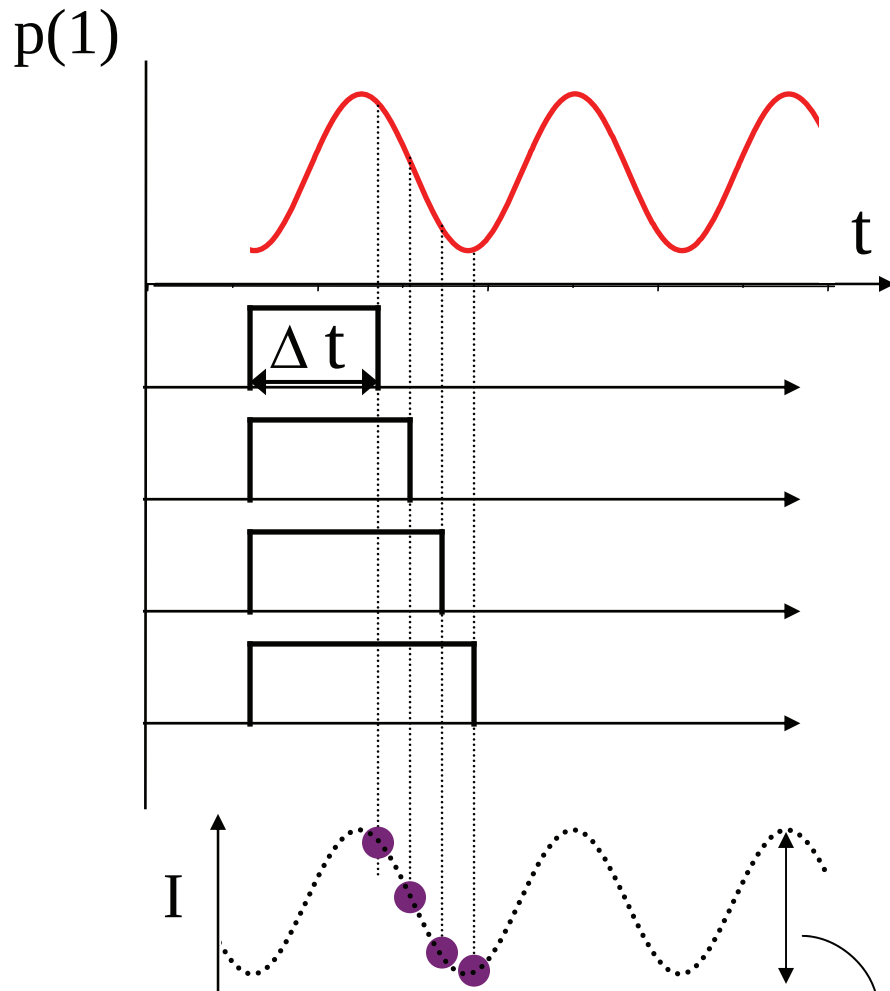
- detect the state $|1\rangle$
- initialize the system to $|0\rangle$

$$E_J \gg \hbar\Gamma_{qp1}$$

Josephson - quasiparticle (JQP) cycle



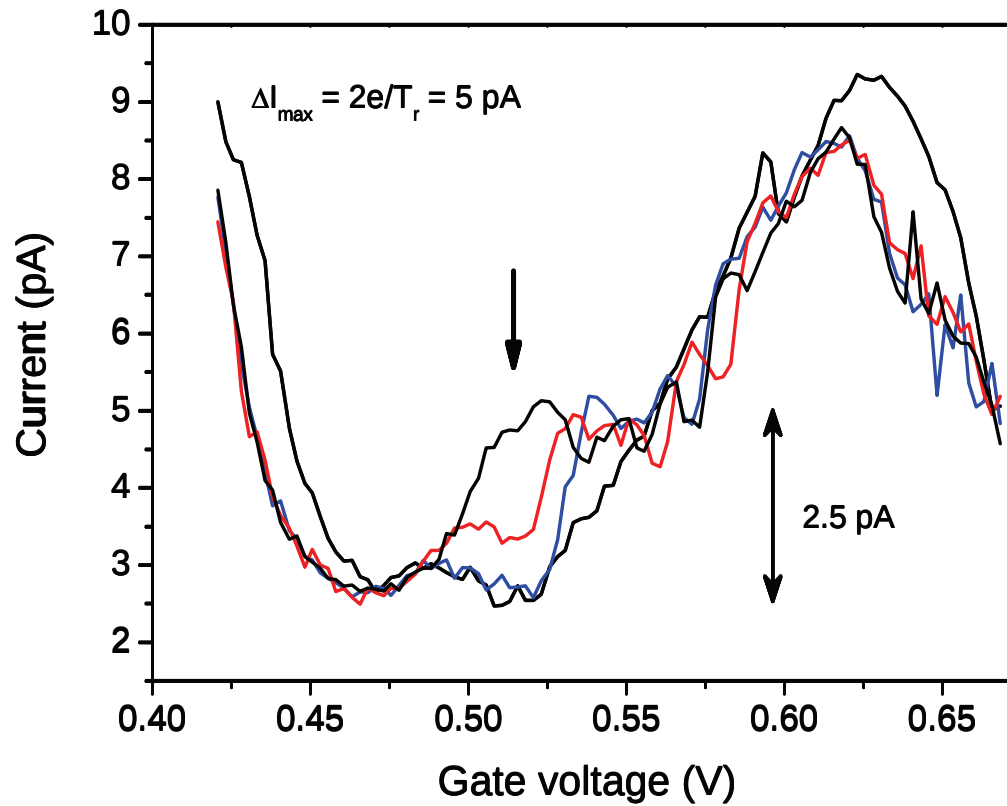
Tracing quantum oscillations: sampling technique



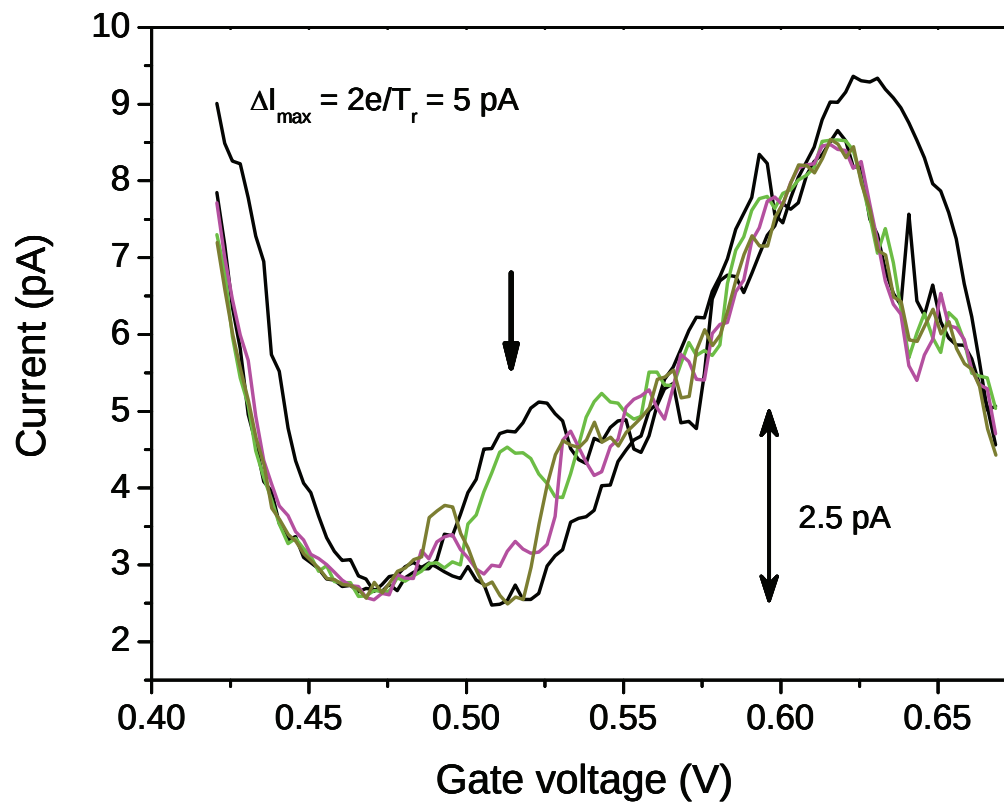
measurement time: 20 ms
pulse array: $20\text{ms}/64 \text{ ns} = 3 \times 10^5$ pulses

$$I = 2e/T_r = 5 \text{ pA}$$

Dc sweep + pulses (1)



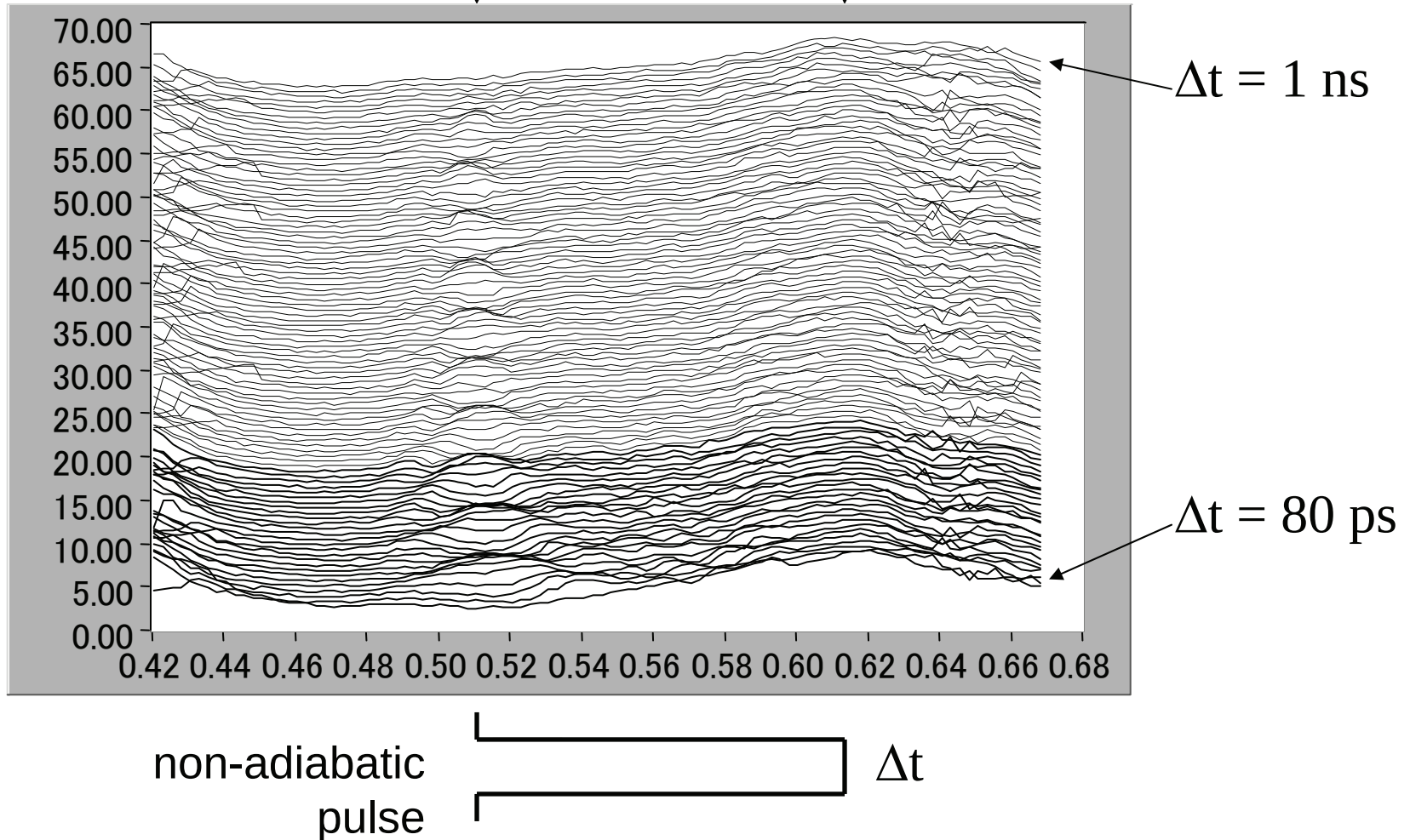
Dc sweep + pulses (2)



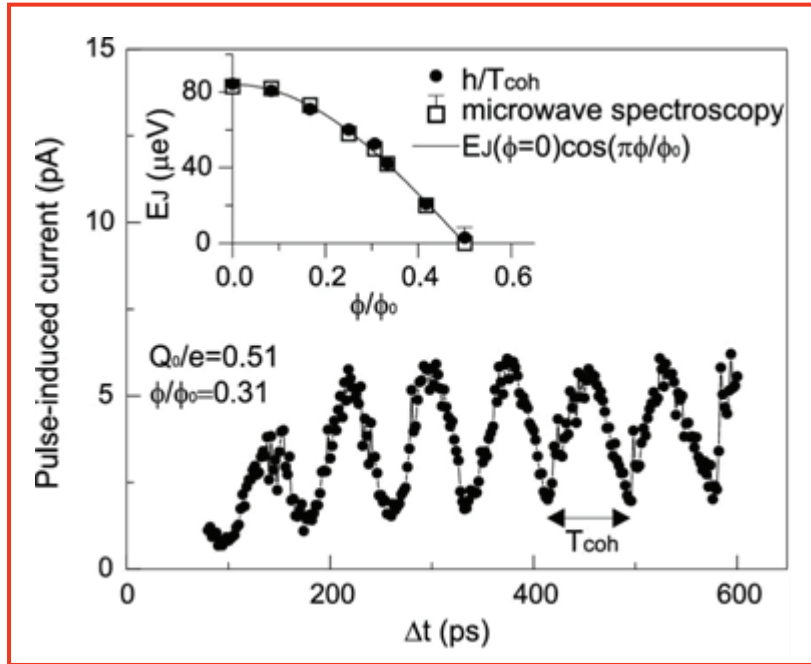
Dc sweep + pulses (4)

pulse-induced
oscillations

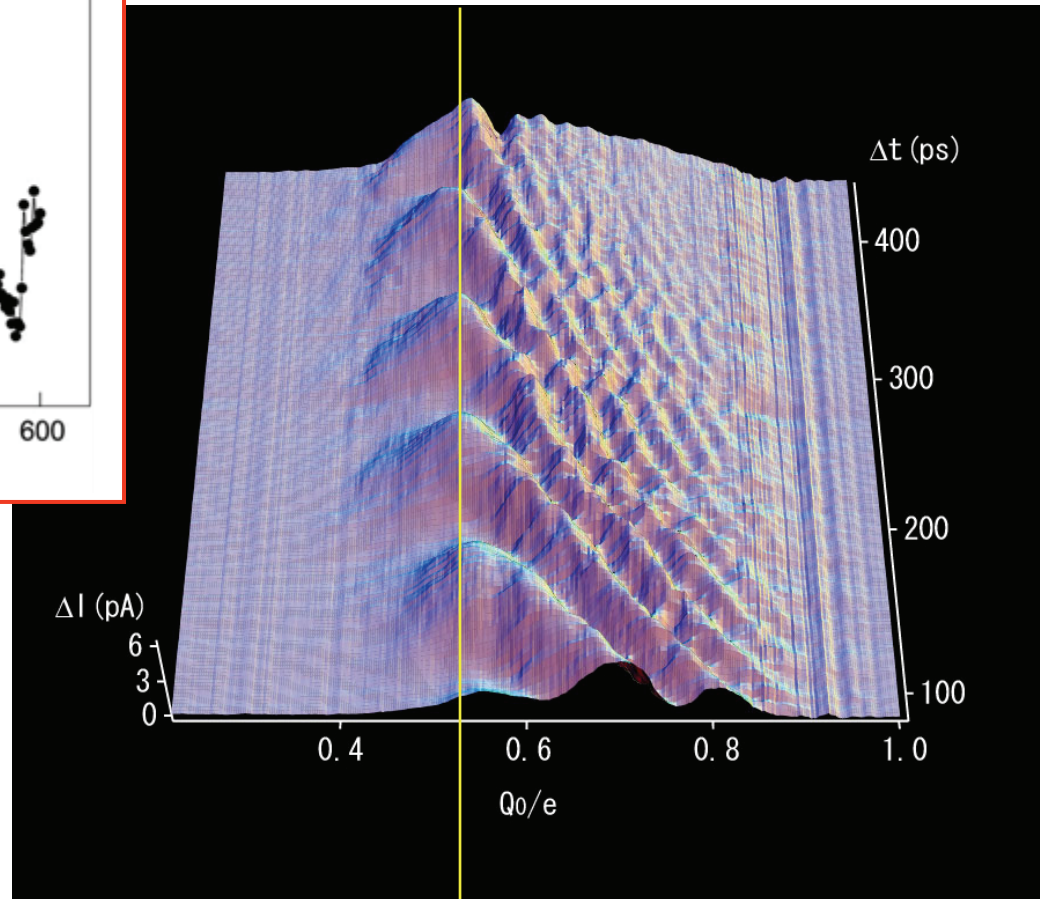
dc JQP peak



Coherent oscillations

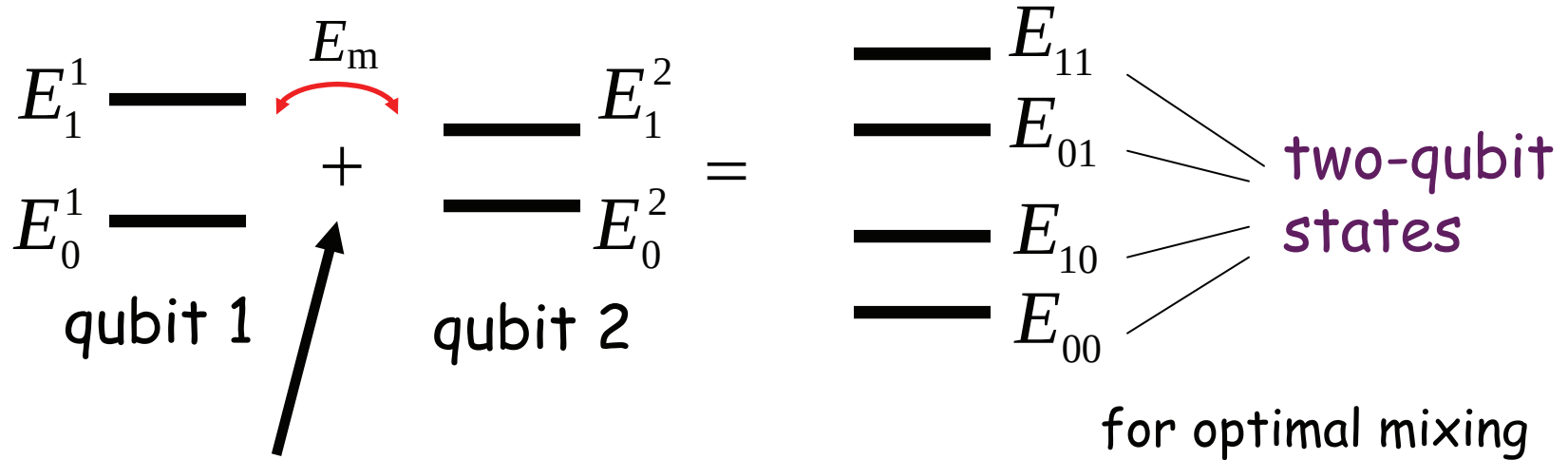


Coherent control
of quantum state



Y. Nakamura *et al.*, Nature **398**, 786 (1999)

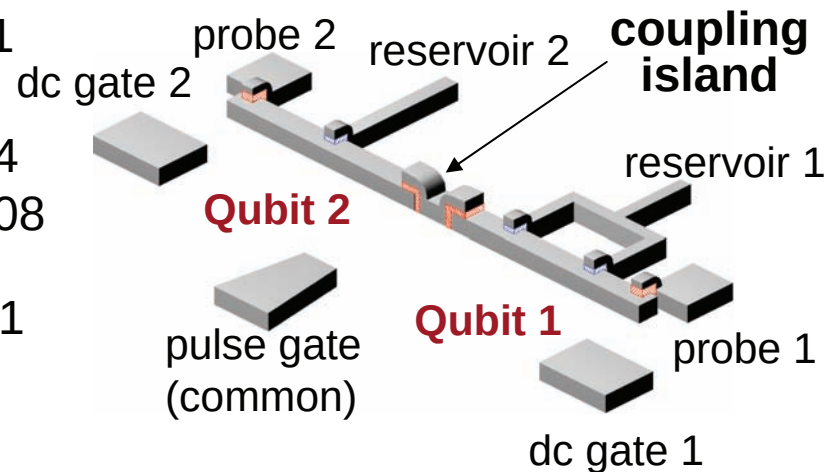
Coupling two charge qubits



electrostatic ← our work

inductive - Yu. Makhlin et al. RMP 2001
 J. Q. You et al. PRL 2002
 J. Lantz et al. cond-mat 2004
 T. Yamamoto et al. PRB 2008

Josephson - F. Plastina et al. PRB 2001



Hamiltonian

charge basis

$$H = \begin{matrix} & \begin{matrix} |00\rangle & |10\rangle & |01\rangle & |11\rangle \end{matrix} \\ \begin{matrix} |00\rangle \\ |10\rangle \\ |01\rangle \\ |11\rangle \end{matrix} & \begin{bmatrix} E_{00} & -\frac{1}{2}E_{J1} & -\frac{1}{2}E_{J2} & 0 \\ -\frac{1}{2}E_{J1} & E_{10} & 0 & -\frac{1}{2}E_{J2} \\ -\frac{1}{2}E_{J2} & 0 & E_{01} & -\frac{1}{2}E_{J1} \\ 0 & -\frac{1}{2}E_{J2} & -\frac{1}{2}E_{J1} & E_{11} \end{bmatrix} \end{matrix}$$

$$E_{J1,2} \sim E_m < E_{c1,2}$$

initial state $|00\rangle$

$$E_{c1}, E_{c2}, E_m$$

$$E_{J1}, E_{J2}$$

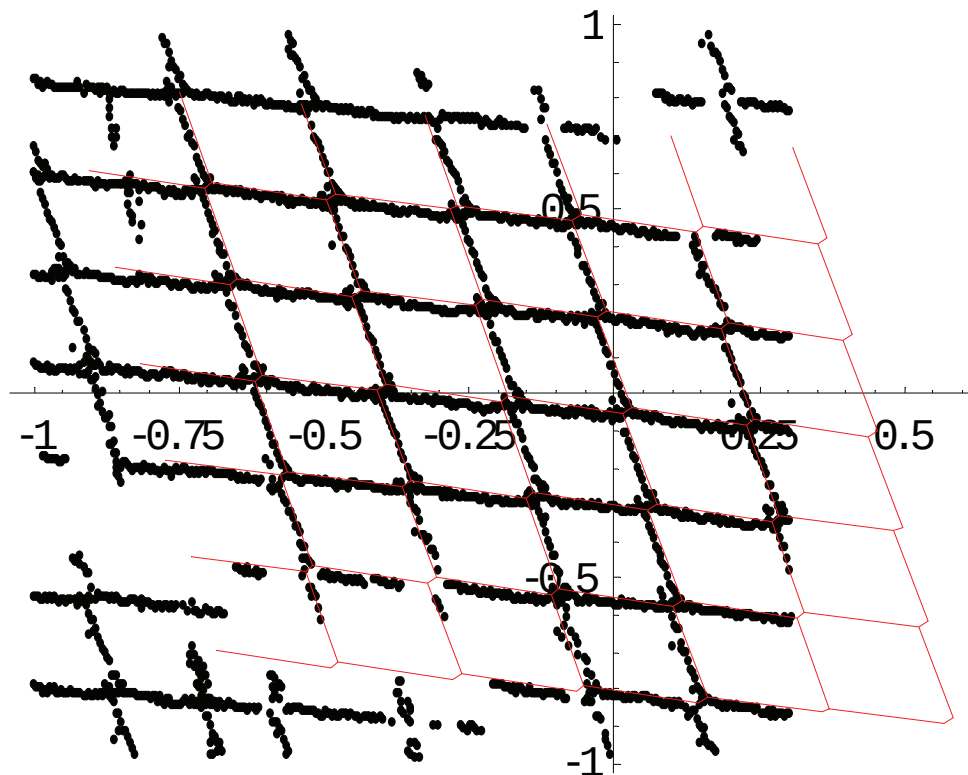
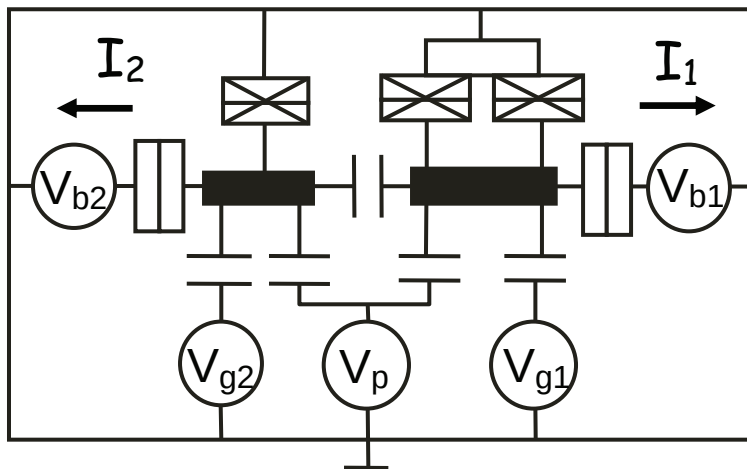
$$E_{n_1 n_2} = E_{c1}(n_{g1} - n_1)^2 + E_{c2}(n_{g2} - n_2)^2 + E_m(n_{g1} - n_1)(n_{g2} - n_2)$$

$$E_{c1,2} = 4e^2 C_{\Sigma 2,1} / 2 (C_{\Sigma 1,2} C_{\Sigma 2,1} - C_m^2) \approx 4e^2 / 2 C_{\Sigma 1,2}$$

$$n_{g1,2} = (C_{g1,2} V_{g1,2} + C_p V_p) / 2e$$

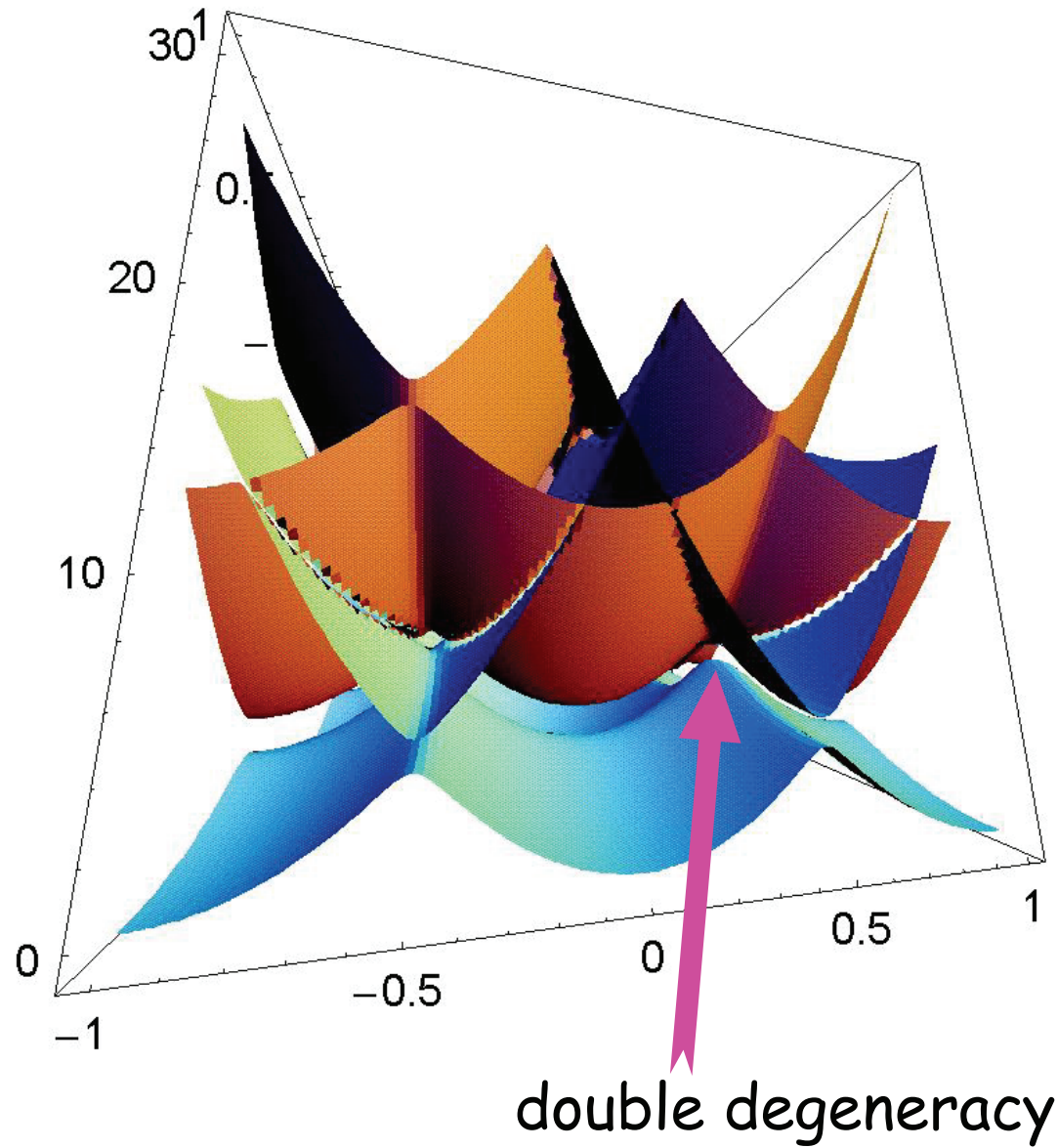
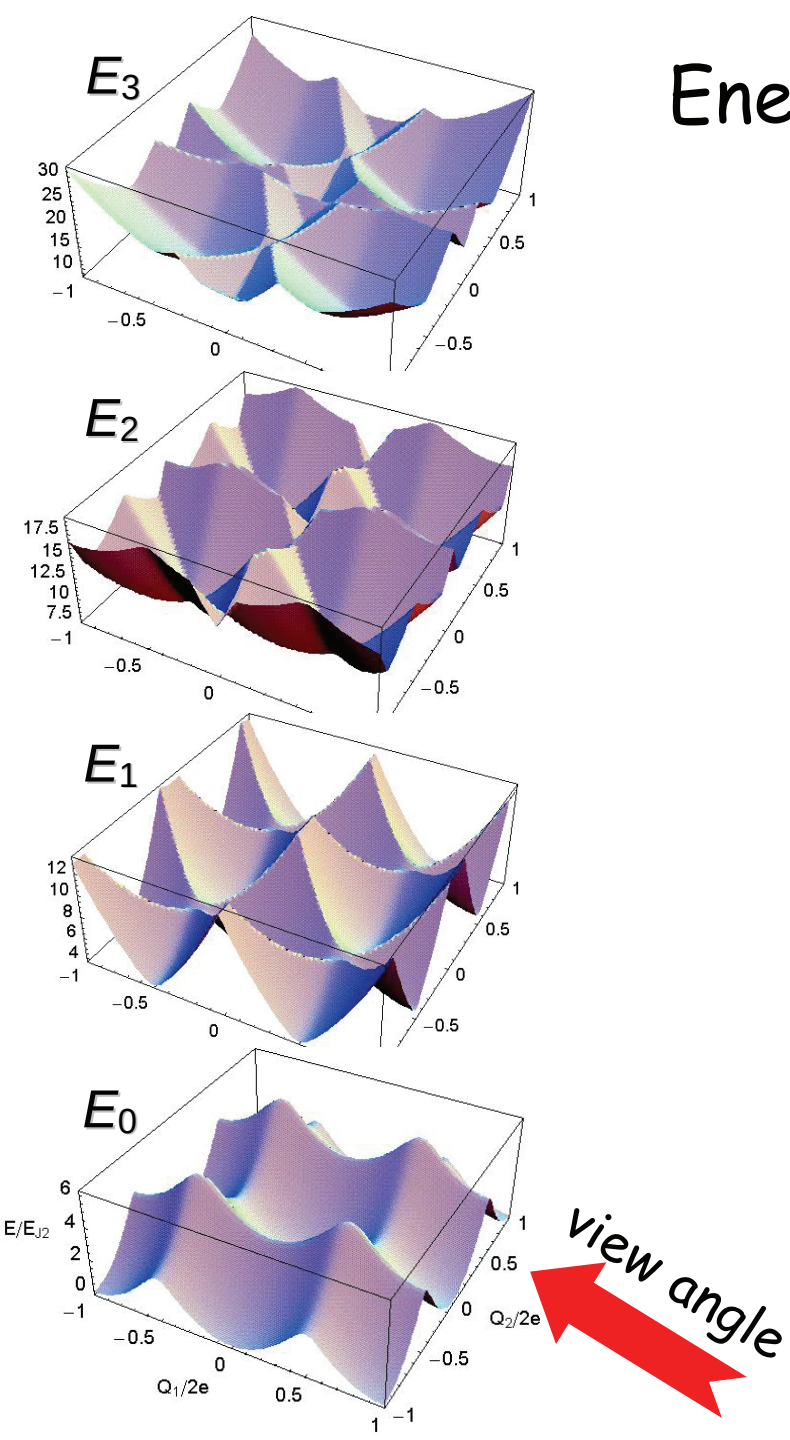
$$E_m = 4e^2 C_m / (C_{\Sigma 1} C_{\Sigma 2} - C_m^2)$$

DC measurements

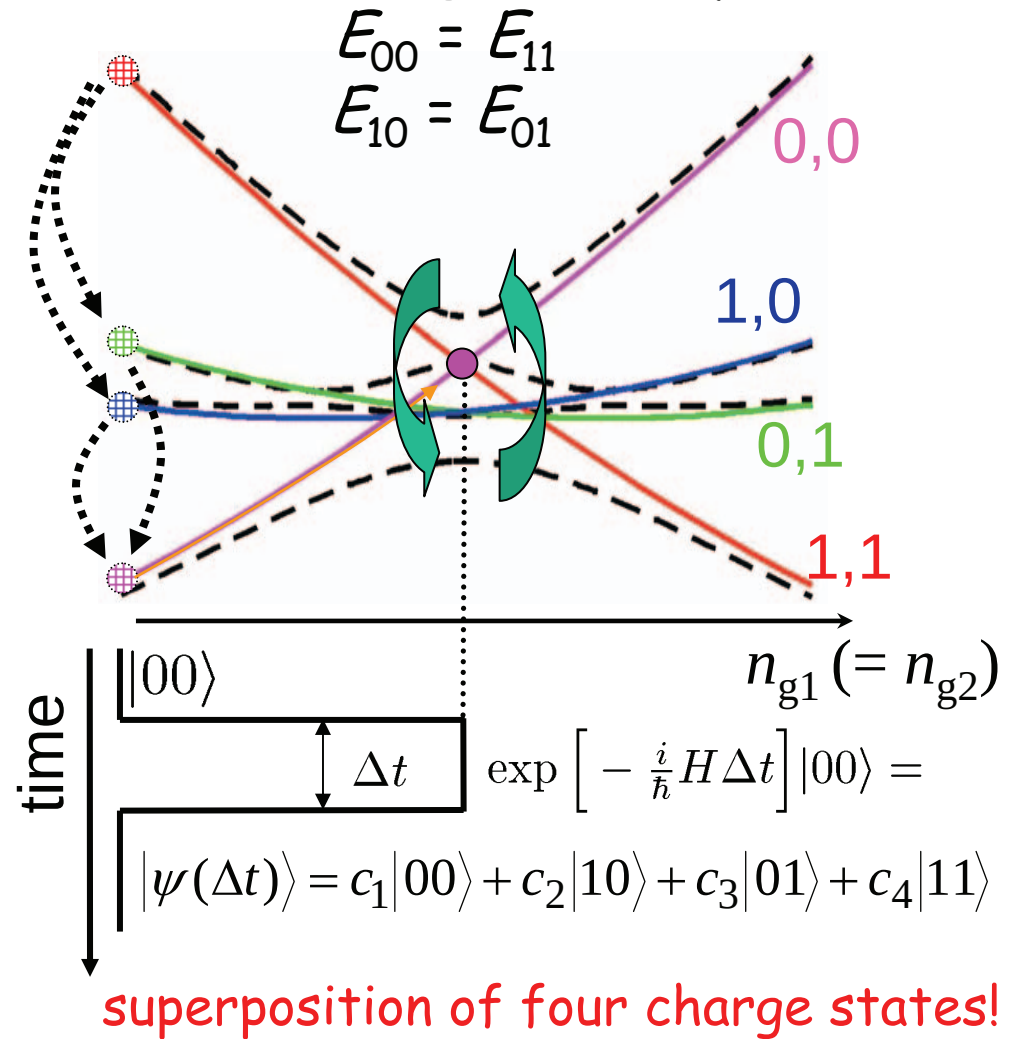
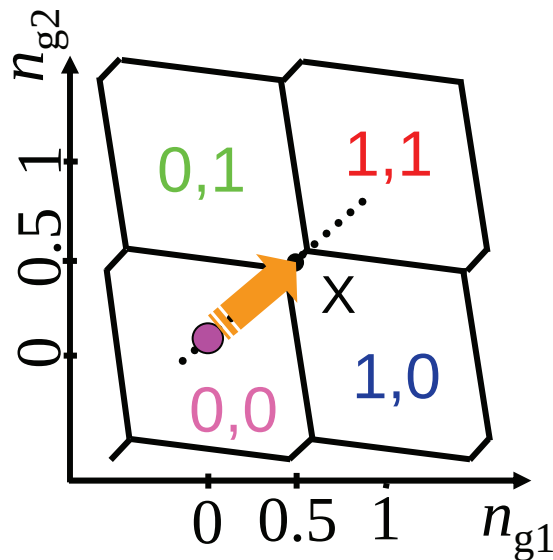
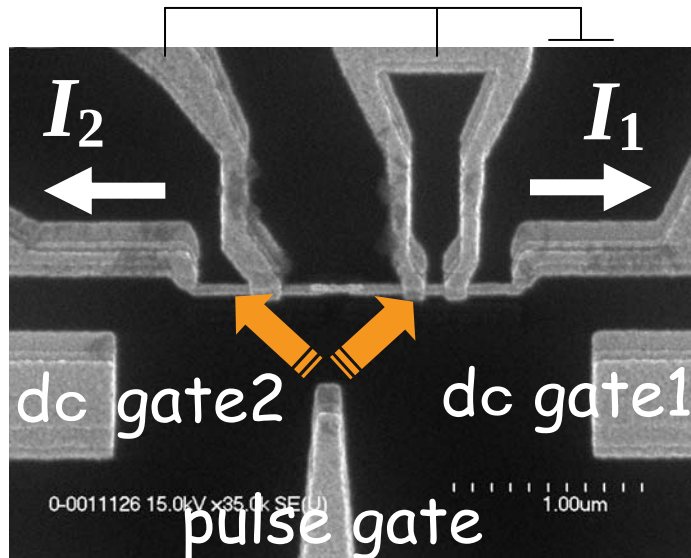


From the fit: $E_{c1} = 120$ GHz
 $E_{c2} = 150$ GHz
 $E_m = 15.7$ GHz

Energy bands



Oscillations at the double degeneracy



Quantum beatings

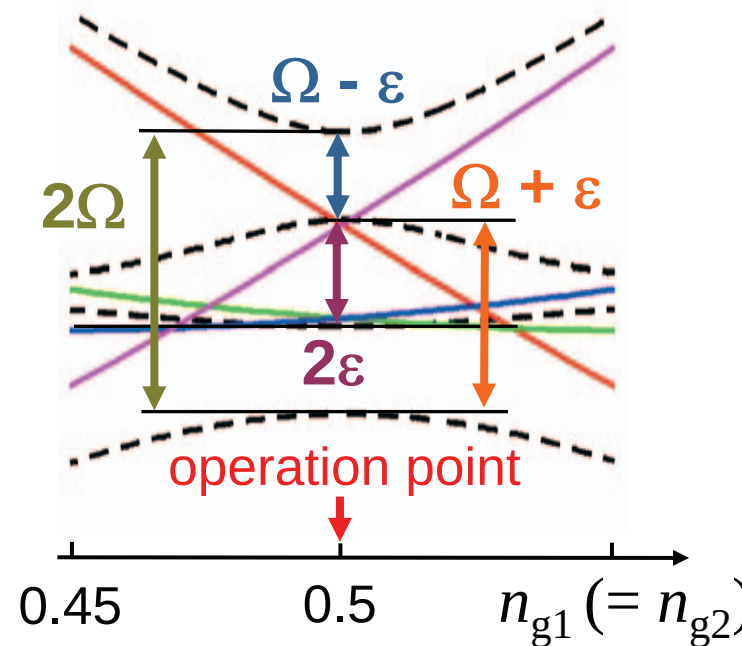
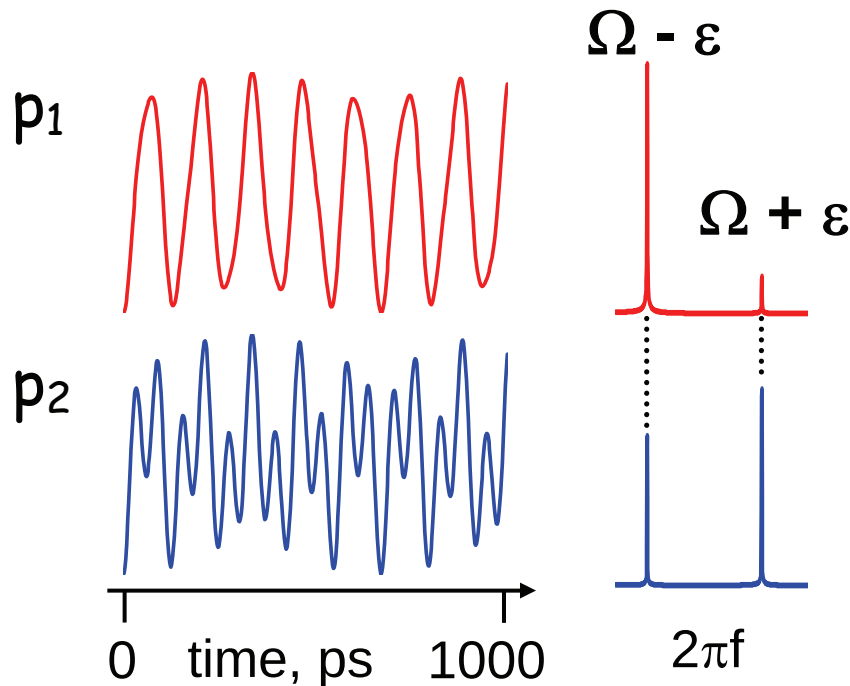
$$|\psi(t)\rangle = \exp\left[-\frac{i}{\hbar} Ht\right] |00\rangle$$

$$|\psi(t)\rangle = c_1|00\rangle + c_2|10\rangle + c_3|01\rangle + c_4|11\rangle$$

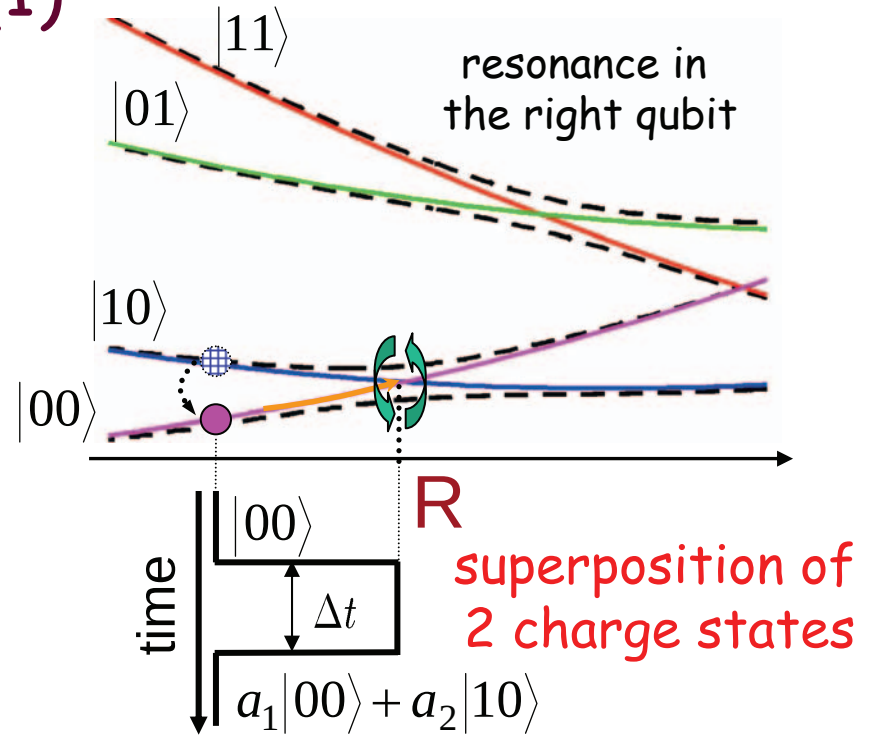
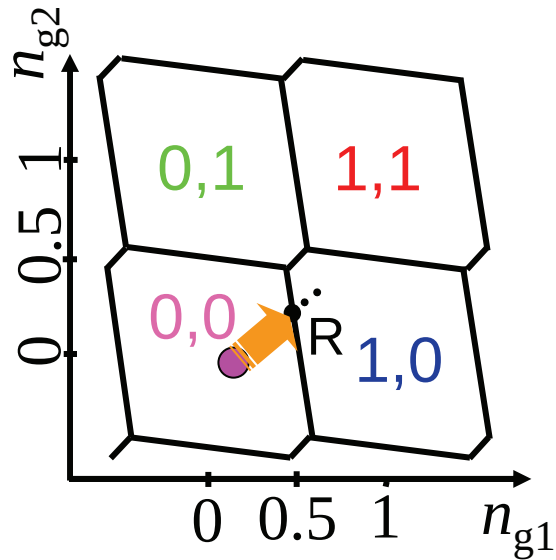
$$I_2 \propto p_2(1) \equiv |c_3|^2 + |c_4|^2 =$$

$$= \frac{1}{4} \left[2 - \underline{(1 - \chi) \cos(\Omega + \varepsilon) \Delta t} - \underline{(1 + \chi) \cos(\Omega - \varepsilon) \Delta t} \right]$$

χ	$= \frac{E_{J1}^2 - E_{J2}^2 + (E_m/4)^2}{4\hbar^2 \Omega \epsilon}$
Ω	$= \sqrt{\Delta^2 + (E_m/4\hbar)^2}$
ϵ	$= \sqrt{\delta^2 + (E_m/4\hbar)^2}$
Δ	$= (E_{J2} + E_{J1})/2\hbar$
δ	$= (E_{J2} - E_{J1})/2\hbar$

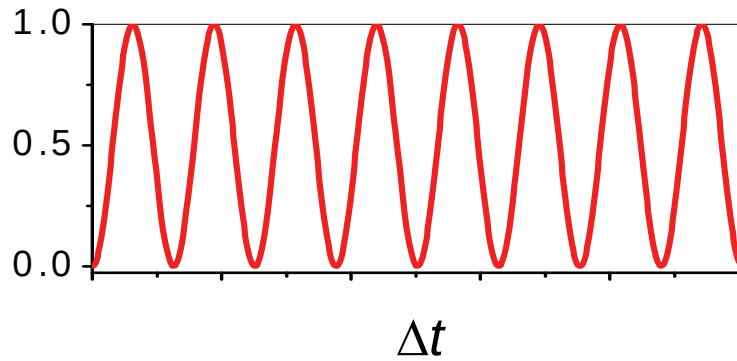


Independent oscillations (1)

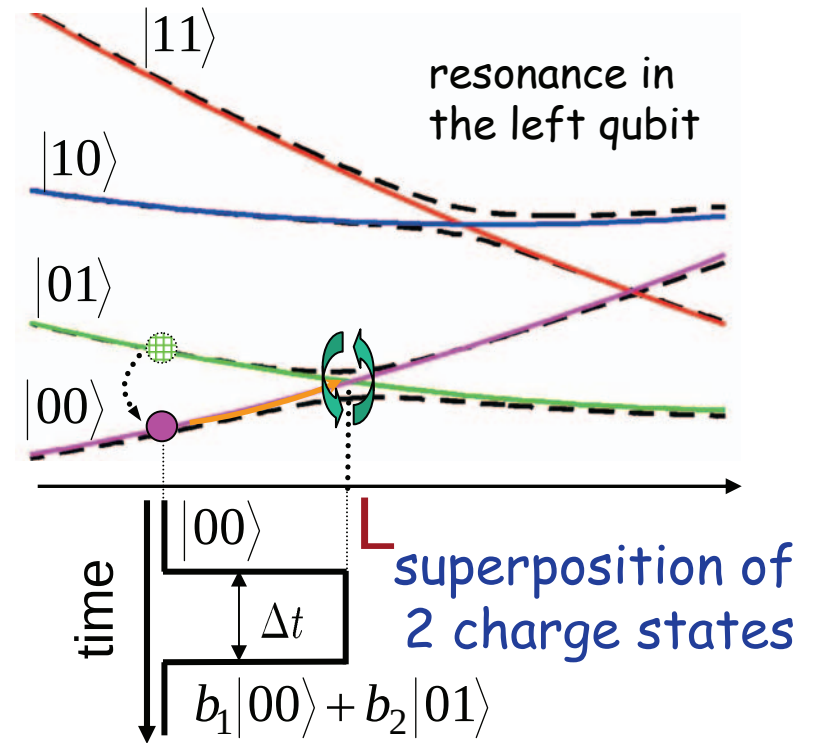
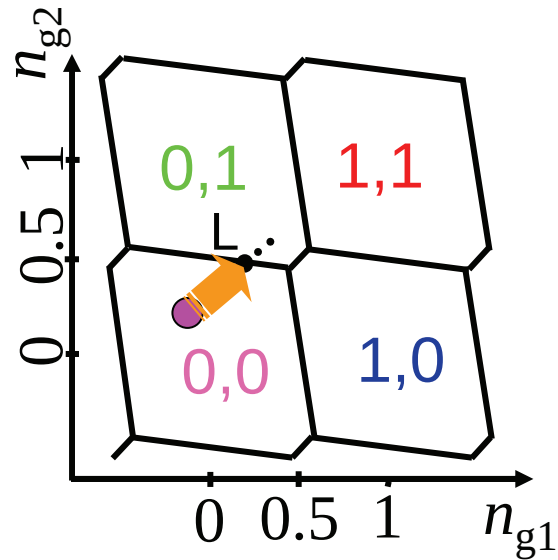


$$I_1 \propto p_1(1) \equiv |a_2|^2 = \frac{1}{2} \left[1 - \cos\left(\frac{E_{J1}}{\hbar} t\right) \right]$$

oscillations in qubit 1

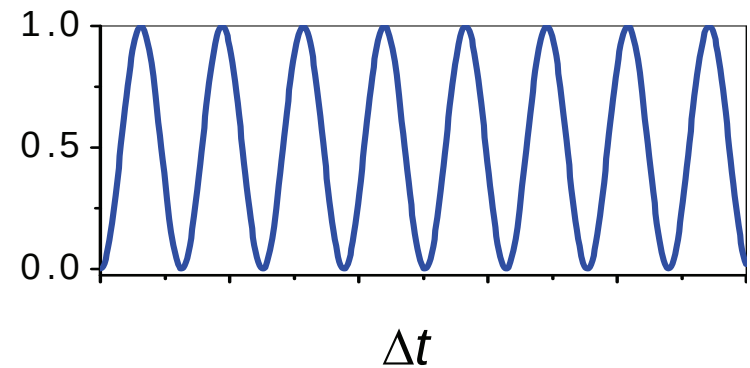


Independent oscillations (2)

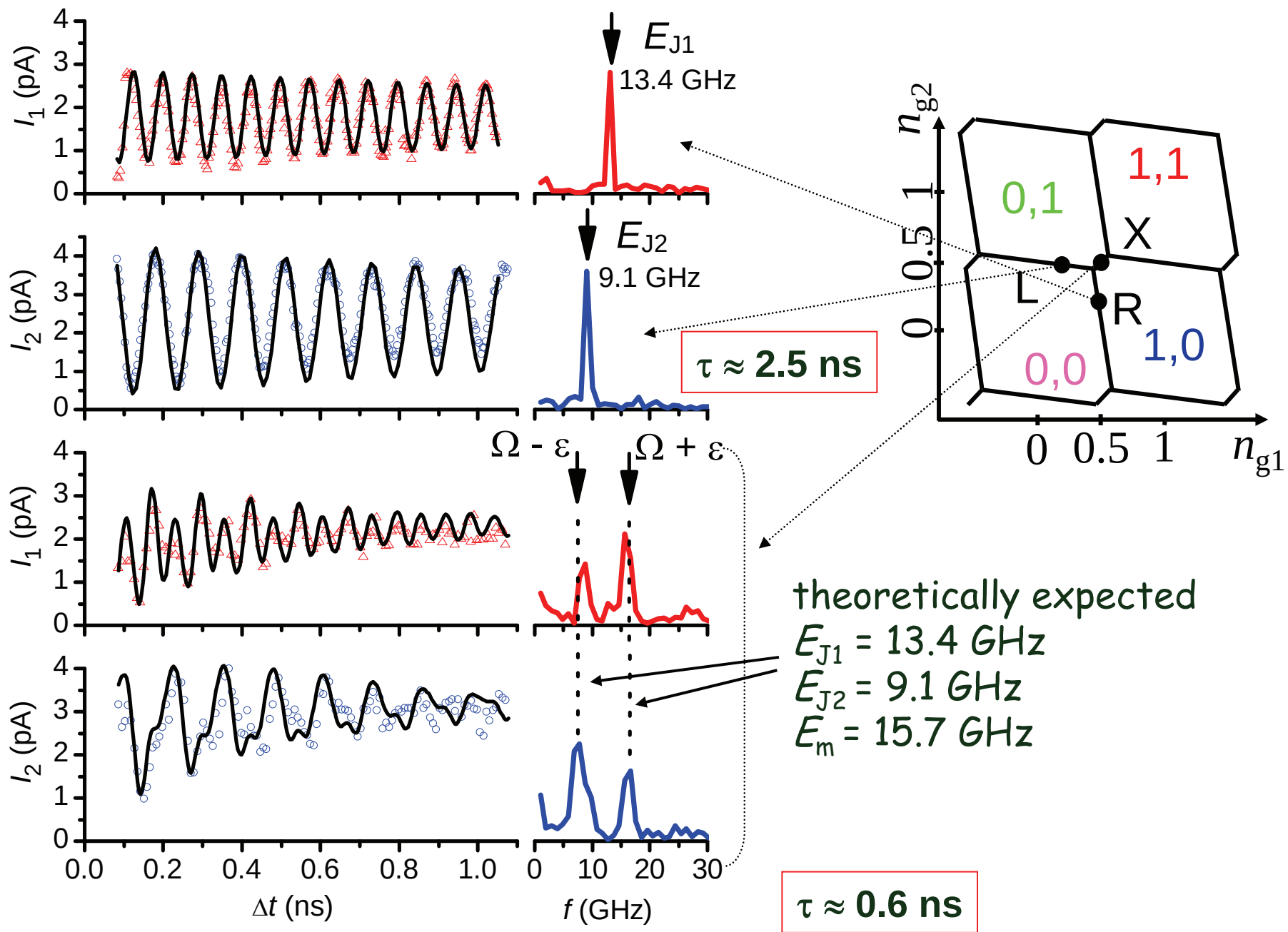


$$I_2 \propto p_2(1) \equiv |b_2|^2 = \frac{1}{2} \left[1 - \cos\left(\frac{E_{J2}}{\hbar} t\right) \right]$$

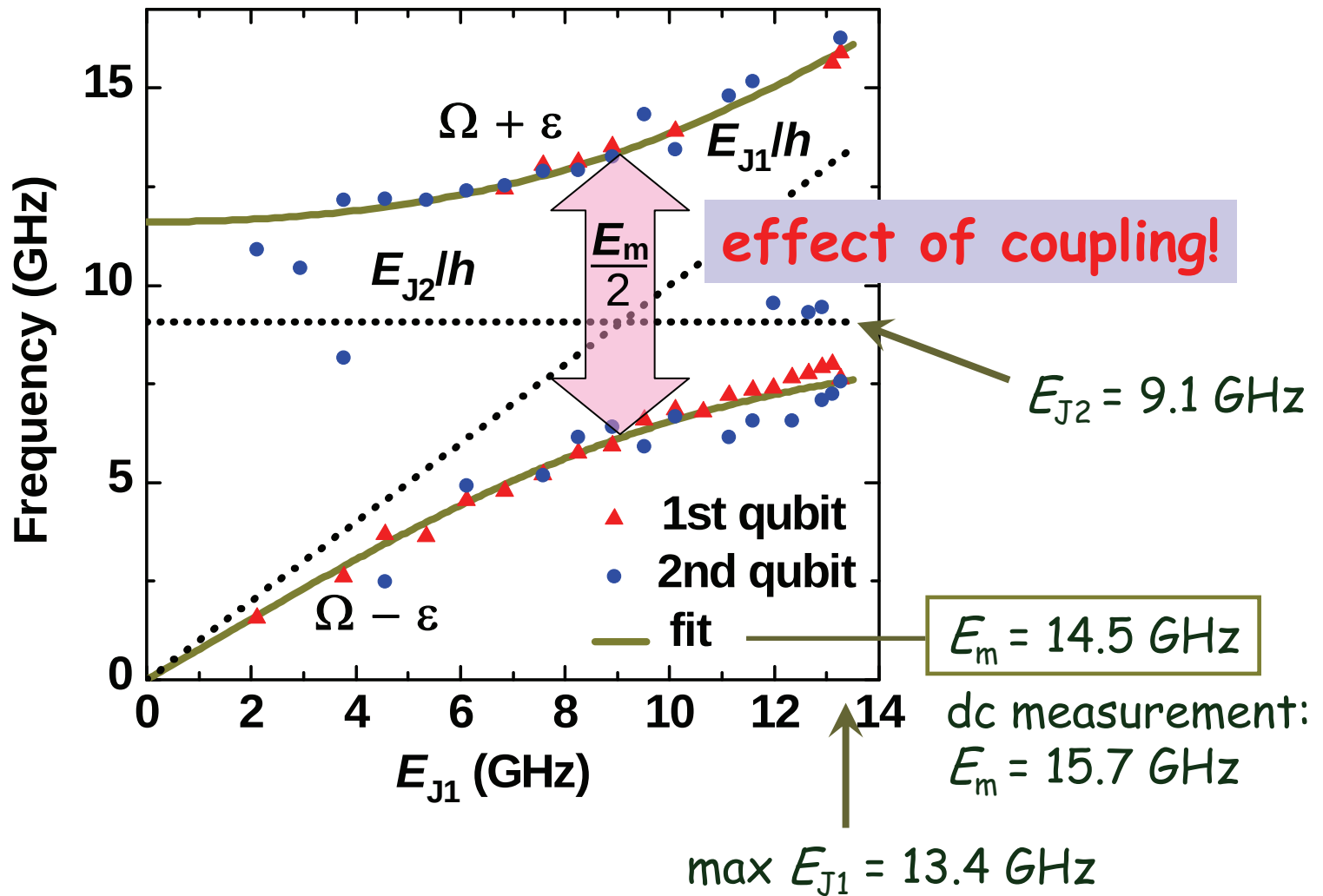
oscillations in qubit 2



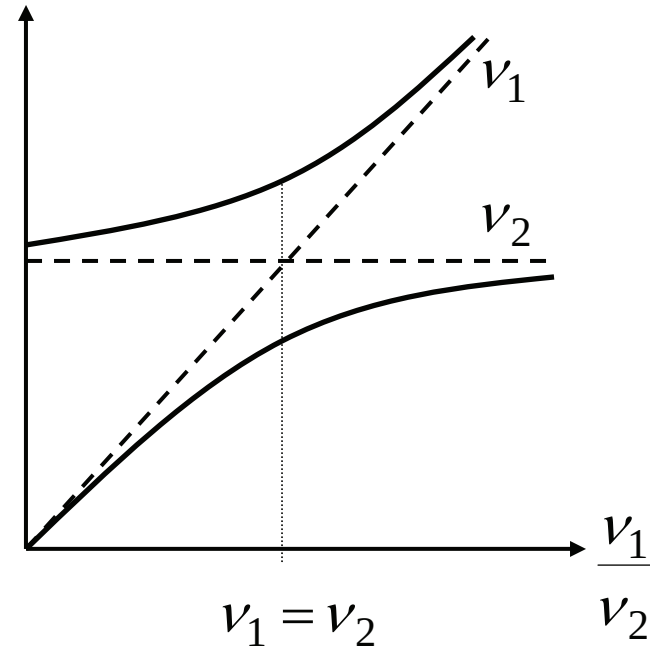
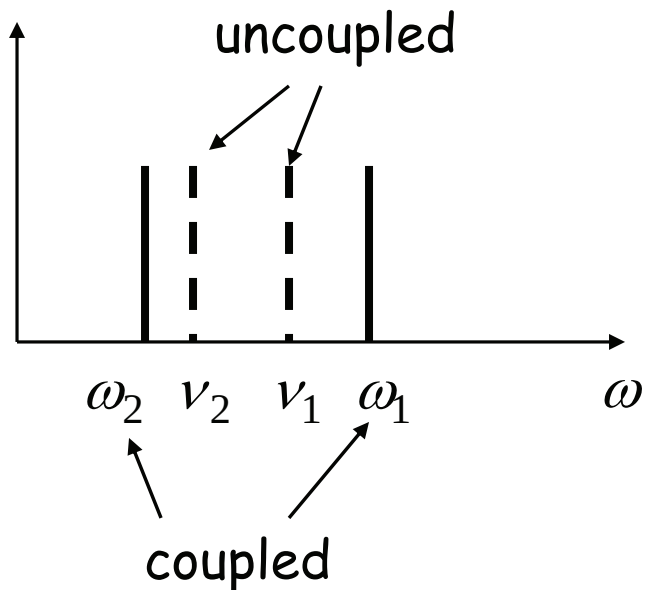
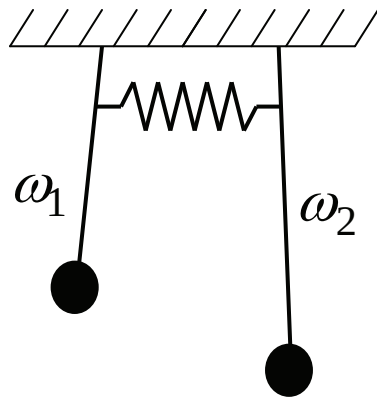
Quantum beatings: experiment



E_{J1} dependence of frequencies



Two coupled classical oscillators



What is the difference?

oscillates:

C.: physical parameter x

Q.: $p(x)$ to be 0 or 1

Entanglement of two coupled qubits

Entangled qubits: $|\psi\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$

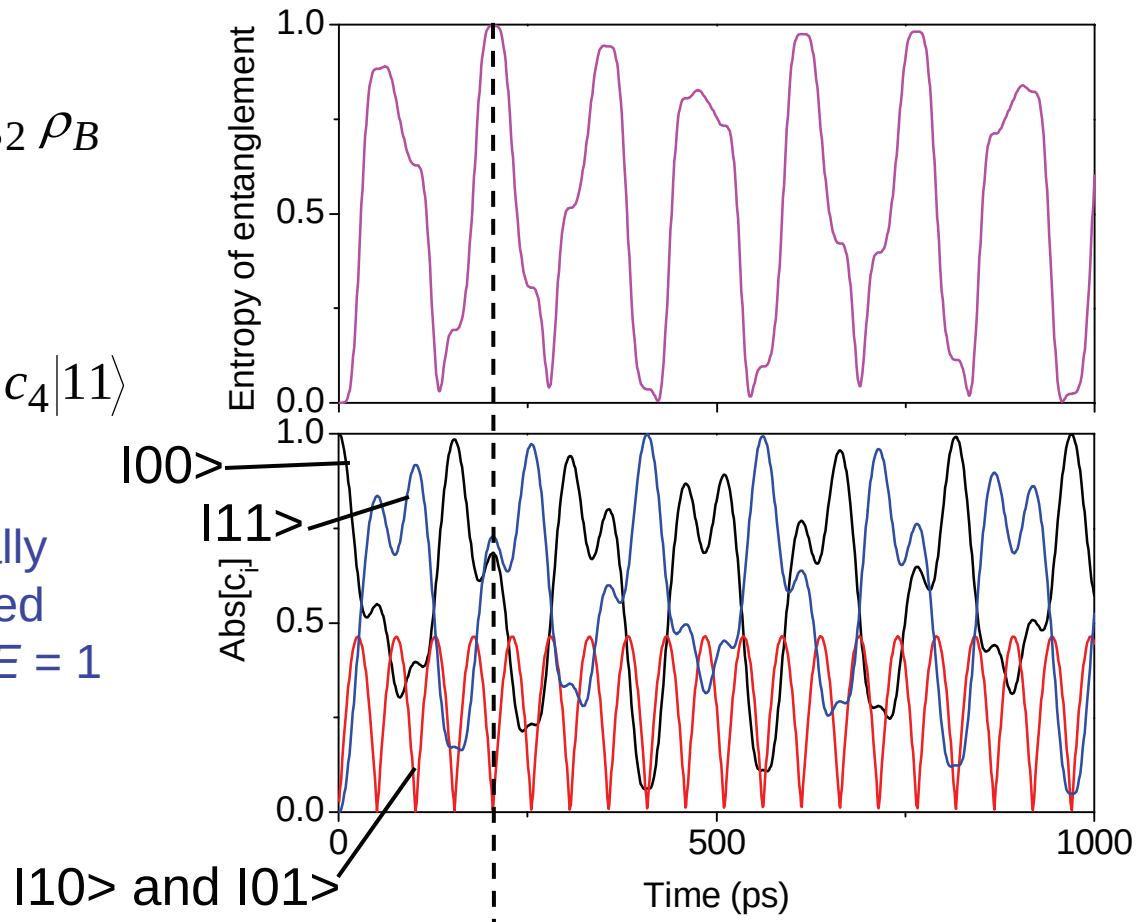
Entropy of entanglement:

$$E = -\text{Tr} \rho_A \log_2 \rho_A = -\text{Tr} \rho_B \log_2 \rho_B$$

our qubits $E_{J1} = 9.1 \text{ GHz}$
 $E_{J2} = 9.1 \text{ GHz}$
 $E_m = 14.5 \text{ GHz}$

If $|\psi\rangle = c_1|00\rangle + c_2|10\rangle + c_3|01\rangle + c_4|11\rangle$

$$\left. \begin{aligned} |\psi\rangle &= \frac{1}{\sqrt{2}}(|00\rangle \pm |11\rangle) \\ |\psi\rangle &= \frac{1}{\sqrt{2}}(|10\rangle \pm |01\rangle) \end{aligned} \right\} \text{maximally entangled states, } E = 1$$



"almost" maximally entangled state

Status of Josephson charge qubits:

stage of proof of concepts passed

- ✓ • first solid-state qubit: coherent oscillations
- ✓ • 2 qubits coupled: beatings, avoided level crossing
- ✓ • conditional gate operation
- ✓ • single-shot readout
- ? • sources of decoherence